

# A Direct Constrained Optimization Method for the Kohn-Sham Equations

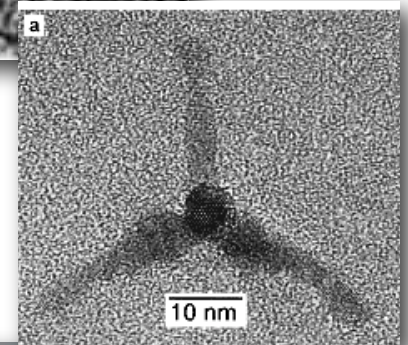
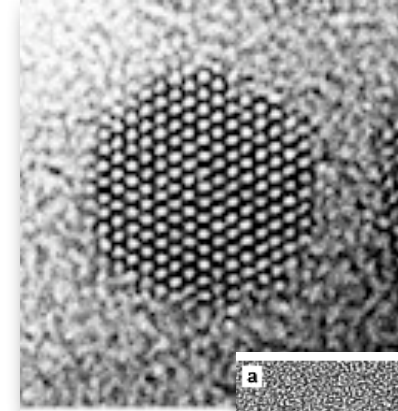


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North Carolina State    November 13, 2008

# What do all of these have in common?



# On using mathematics for chemistry

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*Every attempt to employ mathematical methods in the study of chemical questions must be considered **profoundly irrational** and contrary to the spirit of chemistry. If mathematical analysis should ever hold a prominent place in chemistry – an aberration which is happily **almost impossible** – it would occasion a rapid and widespread degeneration of that science.*

*Auguste Comte, 1830*

# 100 years later – the problem is solved!

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*...in the Schrödinger equation we very nearly have the mathematical foundation for the solution of the whole problem of atomic and molecular structure ...*

***almost***

*...the problem of the many bodies contained in the atom and the molecule cannot be completely solved without a great further development in mathematical technique.*

*G.N. Lewis, J. Chem. Phys. 1, 17 (1933)*

# Many-body Schrödinger equation

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$$H\Psi_i(r_1, r_2, \dots, r_N) = E_i\Psi(r_1, r_2, \dots, r_N)$$

$$H = -\frac{\hbar^2}{2m} \sum_{i=1}^N \nabla_i^2 + \sum_{i=1}^N v(r_i) + \frac{1}{2} \sum_{i \neq j} \frac{e^2}{|r_i - r_j|}$$

- $\Psi_i$  contains all the information needed to study a system
- $|\Psi_i|^2$  probability density of finding electrons at a certain state
- $E_i$  quantized energy
- Computational work goes as  $10^{3N}$ , where N is the number of electrons

# Density Functional Theory

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- The unknown is simple – the electron density,  $\rho$
- Hohenberg-Kohn Theory
  - There is a unique mapping between the ground state energy and density
  - Exact form of the functional is unknown
- Independent particle model
  - Electrons move independently in an average effective potential field
  - Add correction for correlation
- Good compromise between accuracy and feasibility

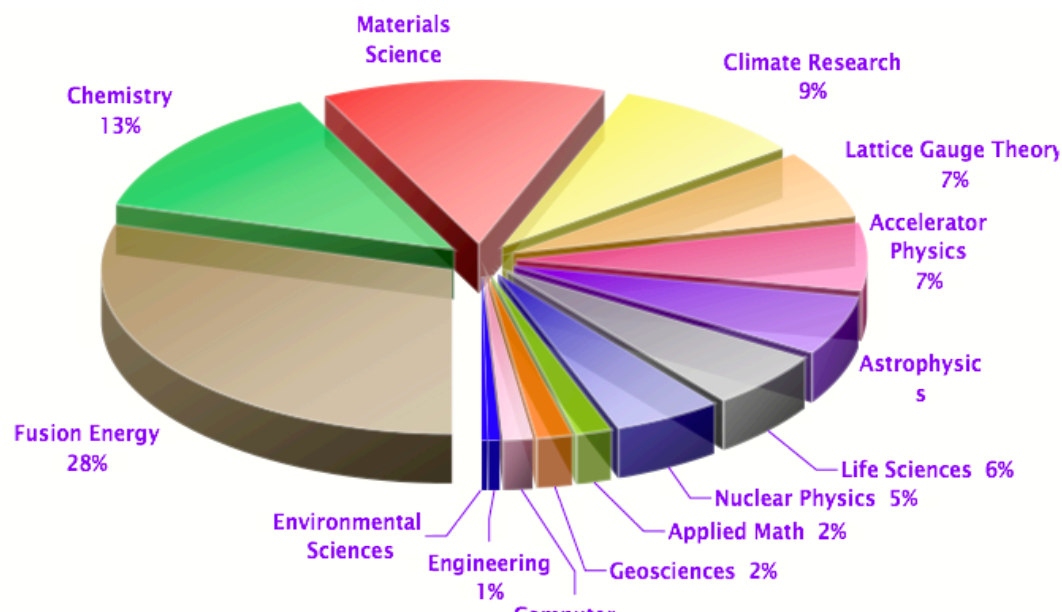


# DFT codes play a major role in computational science

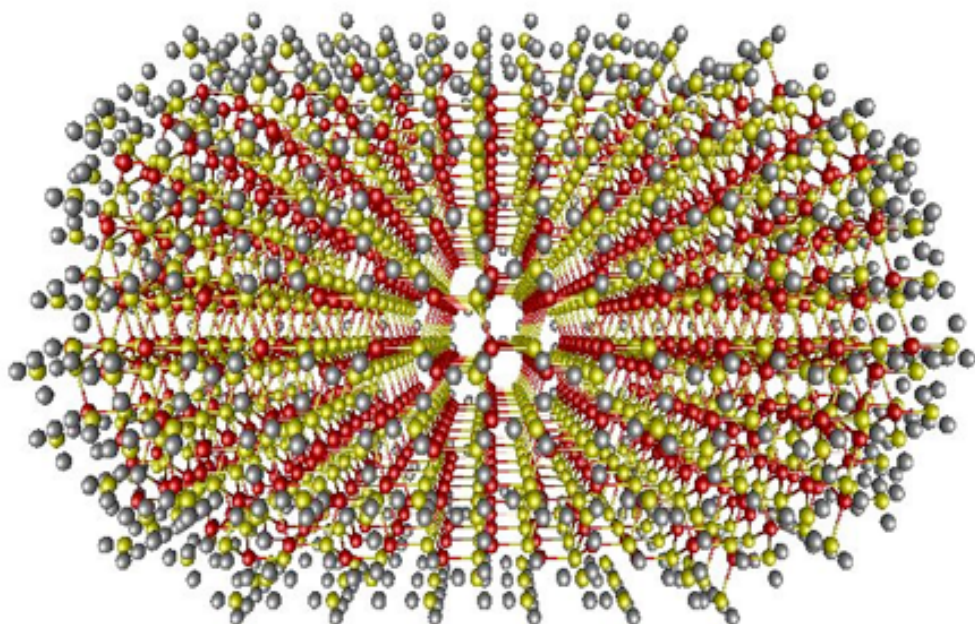


Franklin (NERSC-5): Cray XT4

- 9,660 nodes; 38,640 2.3 GHz cores
- 356 Tflops/s peak
- DFT methods account for 75% of the materials sciences simulations at NERSC, totaling over 5 million hours of computer time in 2006



# Fast forward to today: we can now simulate realistic nanosystems



Advances in ***density functional theory coupled with multinode computational clusters*** now enable accurate simulation of the behavior of multi-thousand atom complexes that mediate the electronic and ionic transfers of solar energy conversion. These new and emerging nanoscience capabilities bring a ***fundamental understanding of the atomic and molecular processes of solar energy utilization within reach.***

The calculated dipole moment of a 2633 atom CdSe quantum rod,  $\text{Cd}_{961}\text{Se}_{724}\text{H}_{948}$ . Using 2560 processors at NERSC the calculation took about 30 hours.

Wang, Zhao, Meza, Phys. Rev. B, 77, 165113 (2008)

*Basic Research Needs for Solar Energy Utilization,*

*Report of the BES Workshop on Solar Energy Utilization, April 18-21, 2005*



# Brief Review of Fundamental Equations

# Kohn-Sham Formulation

- Use  $N$  noninteracting electrons as a reference
- Replace many-particle wavefunctions,  $\Psi_i$ , with single-particle wavefunctions,  $\psi_i$
- Write Kohn-Sham total energy as:

$$E_{total}[\{\psi_i\}] = \frac{1}{2} \sum_{i=1}^{n_e} \int_{\Omega} |\nabla \psi_i|^2 + \int_{\Omega} V_{ion} \rho \\ + \frac{1}{2} \int_{\Omega} \frac{\rho(r)\rho(r')}{|r - r'|} dr dr' + E_{xc}(\rho),$$

$$\rho(r) = \sum_{i=1}^{n_e} |\psi_i(r)|^2, \int_{\Omega} \psi_i \psi_j = \delta_{i,j}, n_e$$

- Exchange-correlation term,  $E_{xc}$ , contains quantum mechanical contributions, plus part of K.E. not covered by first term when using single-particle wavefunctions

# Kohn-Sham Equations

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- Goal is to find the ground state energy by minimizing total energy,  $E_{total}$
- Leads to:

$$H\psi_i = \epsilon_i\psi_i, \quad i = 1, 2, \dots, n_e$$
$$H = \left[ -\frac{1}{2}\nabla^2 + V_{ion}(r) + \int \frac{\rho}{|r - r'|} + V_{xc}(\rho) \right]$$

# Discretization Options

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- Finite difference  $\psi'(r_j) \approx [\psi(r_j + h) - \psi(r_j - h)]/h$
- Finite elements

$$\psi(r) \approx \sum_j^n \alpha_j \phi_j(r), \quad \phi_j(r) \text{ functions with local support}$$

- Local orbital method (good for molecules)
  - Choose  $\phi_j(r)$  as Gaussian or other “nice” functions
- Planewave expansion
  - Choose  $\phi_j(r)$  as  $e^{ig_j \cdot r}$

# Finite Dimensional Problem

We want to find the ground state energy

$$\min E[\{\psi_i\}] = \frac{1}{2} \sum_{i=1}^{n_e} \int_{\Omega} |\nabla \psi_i|^2 + \int_{\Omega} V_{ion} \rho + \frac{1}{2} \int_{\Omega} \frac{\rho(r)\rho(r')}{|r-r'|} dr dr' + E_{xc}(\rho)$$

After discretization we have:

$$\min_{s.t. X^* X = I_{n_e}} E_{KS}(X) \equiv E_{kinetic}(X) + E_{ion}(X) + E_{Hartree}(X) + E_{xc}(X),$$

where

$$\begin{matrix} & n_e \\ N & \boxed{X} \end{matrix}$$

$$E_{kinetic} = \frac{1}{2} \text{trace}(X^* L X)$$

$$E_{ionic} = \text{trace}(X V_{ion} X^*)$$

$$E_{Hartree} = \frac{1}{2} \rho(X)^T L^\dagger \rho(X)$$

$$E_{xc} = \rho(X)^T (\mu_{xc}[\rho(X)])$$

$$\rho(X) = \text{diag}(X X^*)$$



# KKT Conditions

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- KKT conditions

$$\begin{aligned}\nabla_X \mathcal{L}(X, \Lambda) &= 0, \\ X^* X &= I_{n_e}.\end{aligned}$$

- Leads to ***Discretized Kohn-Sham equations***

$$\begin{aligned}H(X)X &= X\Lambda, \\ X^* X &= I_{n_e}, \\ H(X) &= \frac{1}{2}L + V_{ion} + \text{Diag} (L^\dagger \rho(X)) + \text{Diag} g_{xc}(\rho(X))\end{aligned}$$

# Approaches for solving the Kohn-Sham Equations

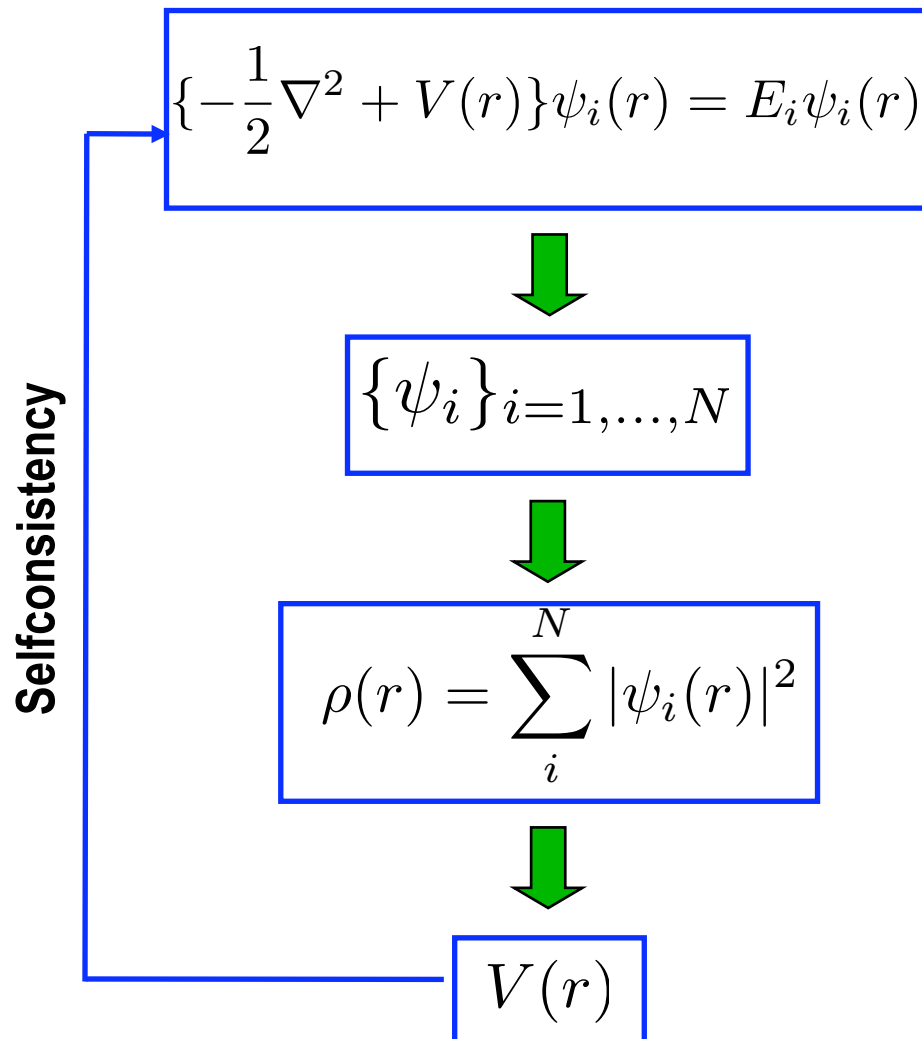
# Solving the Kohn-Sham equations

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- Self-Consistent Field (SCF) iteration
  - view as a linear eigenvalue problem
  - need to precondition
  - usually used with other acceleration techniques to improve convergence
  - no good convergence theory
- Direct Constrained Minimization
  - minimize the total energy directly
  - pose as a constrained optimization problem
  - also requires globalization techniques
- Invariance property

$$\begin{aligned} E(XQ) &= E(X) \\ H(XQ) &= H(X) \end{aligned} \quad \text{for any } Q^*Q = I_{n_e}$$

# Basic SCF Iteration



- Overall Complexity  $O(N^3)$
- Major computational work (for plane wave codes):
  - 3D FFT
  - Orthogonalization
  - Nonlocal potential
- Parallel efficiencies can be quite high
- May converge slowly and sometimes doesn't converge at all
- Energy need not decrease monotonically

# When can we expect SCF to work?

- SCF seeks to minimize a sequence of surrogate models
- Gradients match at  $x^{(i)}$ , i.e.  $\nabla E(x^{(i)}) = \nabla E_{sur}(x^{(i)})$
- Consider simple 2D example:

$$E(x) = \frac{1}{2}x^T Lx + \frac{\alpha}{4}\rho(x)^T L^{-1}\rho(x)$$

$$L = \begin{pmatrix} 2 & -1 \\ -1 & 2 \end{pmatrix}, \quad x = \begin{pmatrix} x_1 \\ x_2 \end{pmatrix}, \quad \rho(x) = \begin{pmatrix} x_1^2 \\ x_2^2 \end{pmatrix}$$

$$\begin{array}{ll} \min E(x) \\ \text{s.t. } x_1^2 + x_2^2 = 1 \end{array}$$

$$\left[ L + \alpha \text{Diag}(L^{-1}\rho(x)) \right] x = \lambda_1 x$$



# Convergence of SCF on toy problem

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$\alpha$	$\ res\ $	iter
2	$1.47 \cdot 10^{-11}$	14
3	$4.01 \cdot 10^{-11}$	22
4	$3.78 \cdot 10^{-11}$	35
5	$7.60 \cdot 10^{-11}$	58
6	$1.07 \cdot 10^{-10}$	126
7	$1.41 \cdot 10^{-1}$	200
8	$1.03 \cdot 10^0$	200

Level sets of surrogate

$X^T X = 1$  constraint

Level sets  
of Energy

SCF step is  
too long!



# Improving SCF

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- Construct better surrogate – cannot afford to use local quadratic approximations (Hessian too expensive)
- **Charge mixing** to improve convergence; related to Broyden methods
- Use **trust region** to restrict the update to stay within a neighborhood of the gradient matching point
  - Level-Shifting (Saunders & Hillier 1973)
  - Cances & LeBris 2000
  - TRSCF – Thogersen, Olsen, Yeager & Jorgensen 2004; Francisco, Martinez, Martinez 2006; Yang, Meza, Wang (2007)

# Charge Mixing Schemes

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- Simple mixing

$$\rho^{(i+1)} \leftarrow \tau \rho_{in}^{(i)} + (1 - \tau) \rho_{out}^{(i)}, \quad 0 < \tau < 1.$$

- Pulay mixing (Direct Inversion of Iterative Subspace)

$$\rho^{(i+1)} = \sum_{j=1}^i \alpha_j \rho^{(j)}, \quad \sum_{j=1}^i \alpha_j = 1$$

- Broyden mixing

$$\rho^{(i+1)} = \rho^{(i)} + \tau C_{i+1} r_i$$

- Anderson mixing

$$\rho^{(i+1)} = \rho^{(i)} + \tau r_i + (S_i - \tau Y_i) Y_i^\dagger r_i$$

# Trust Region Subproblem

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- Solve

$$\min \quad E_{sur}(x)$$

s.t.

$$x^T x = 1,$$

$$\|xx^T - x^{(i)}(x^{(i)})^T\|_F^2 \leq \Delta \quad \text{trust region constraint}$$

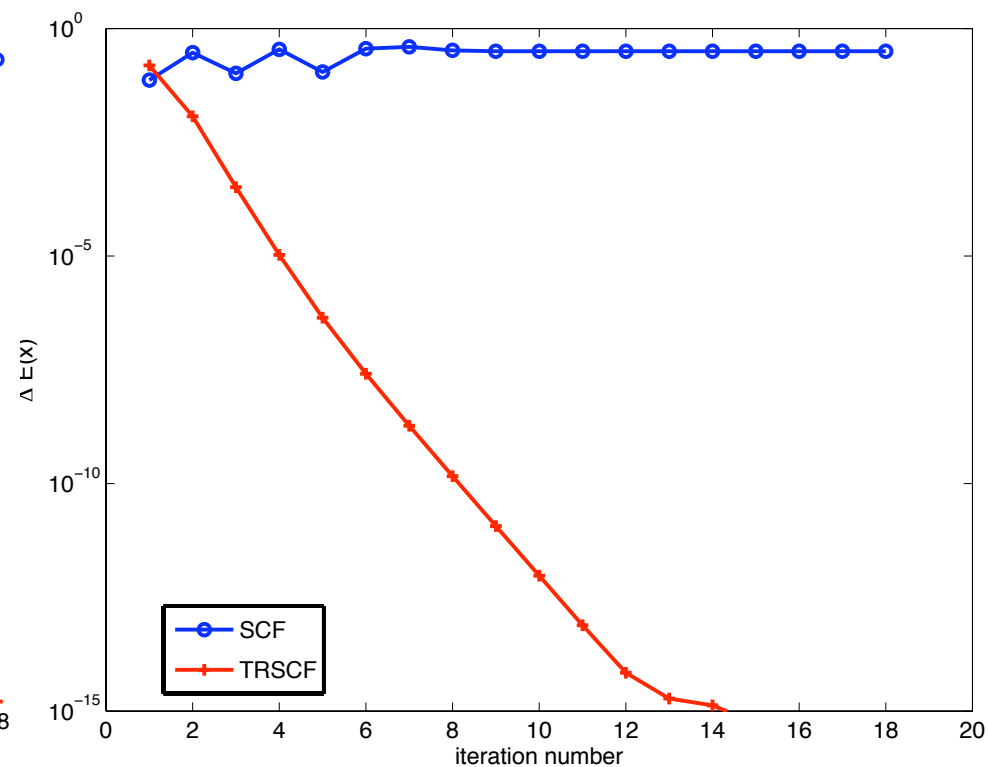
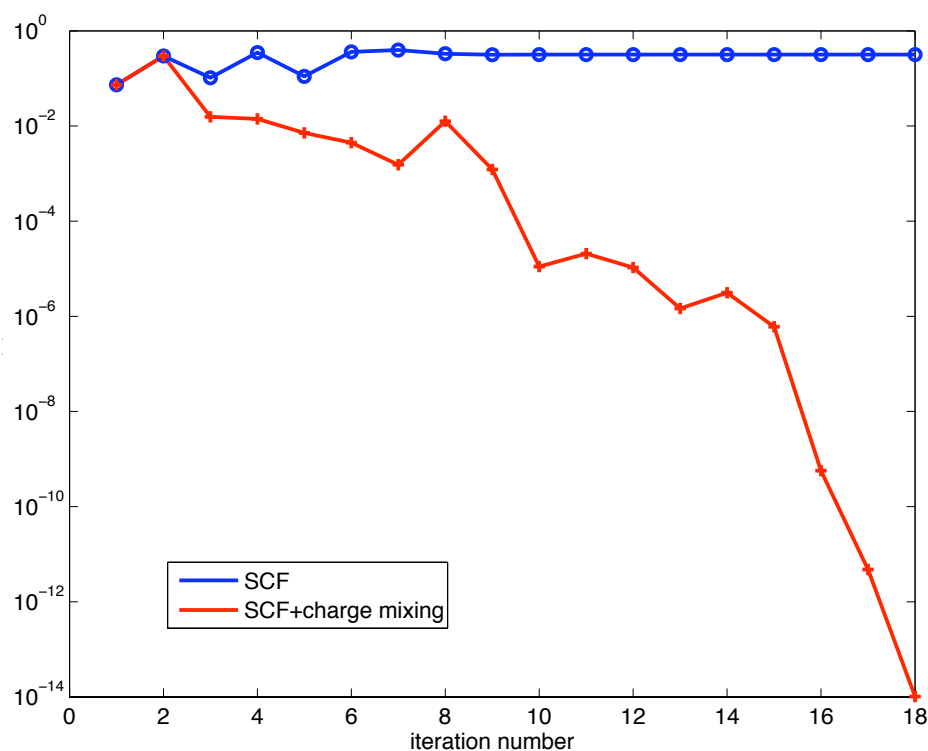
- Equivalent to solving

$$\begin{aligned} \left[ H(x^{(i)}) - \sigma x^{(i)}(x^{(i)})^T \right] x &= \lambda x \\ x^T x &= 1 \end{aligned}$$

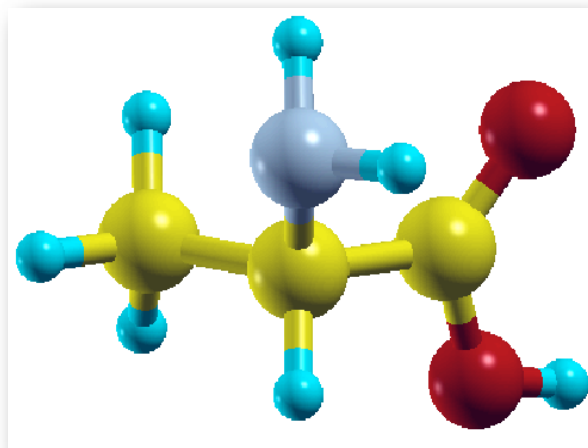
- $\sigma$  is a penalty parameter (Lagrange multiplier for TR)



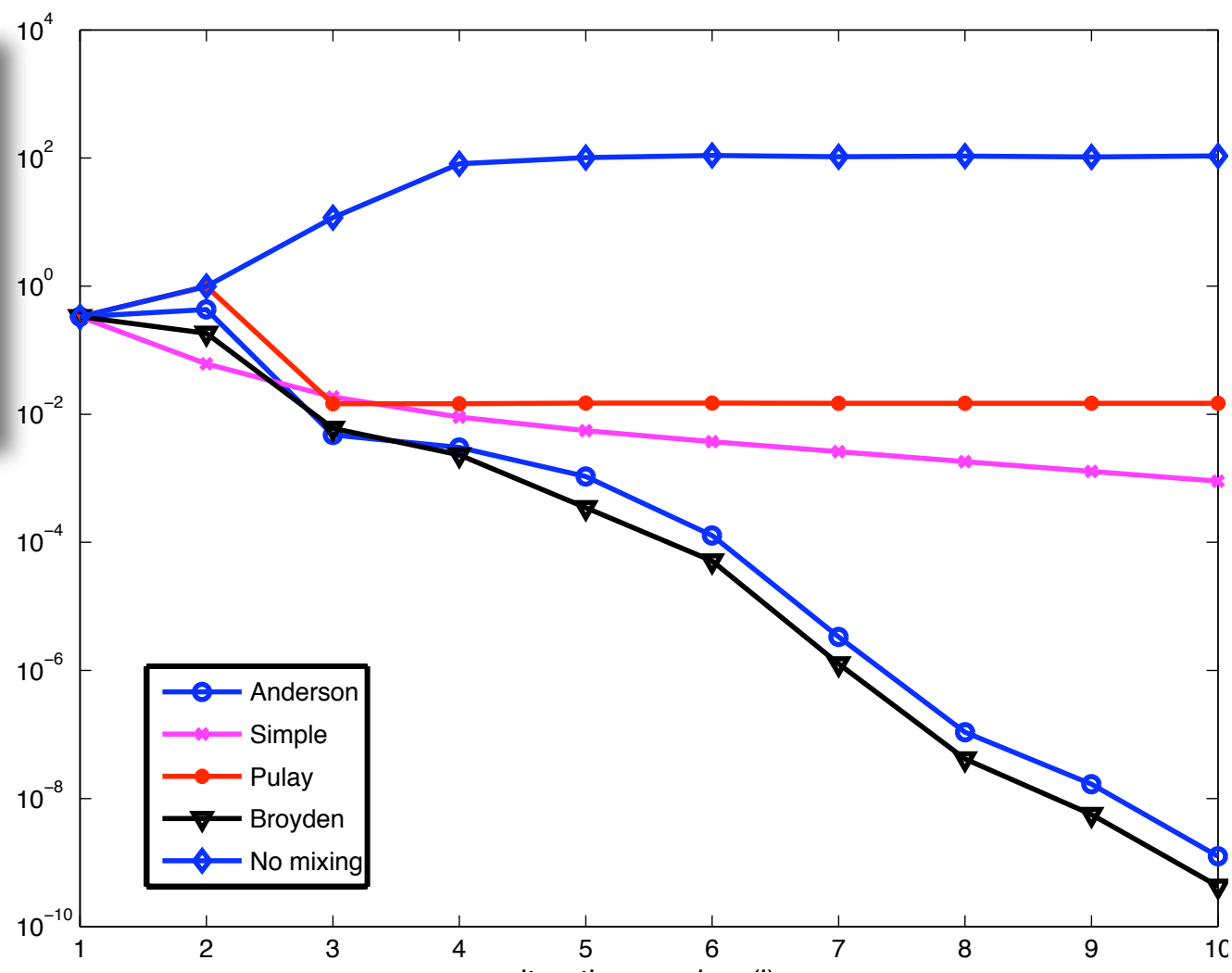
# Comparison of TRSCF vs. mixing



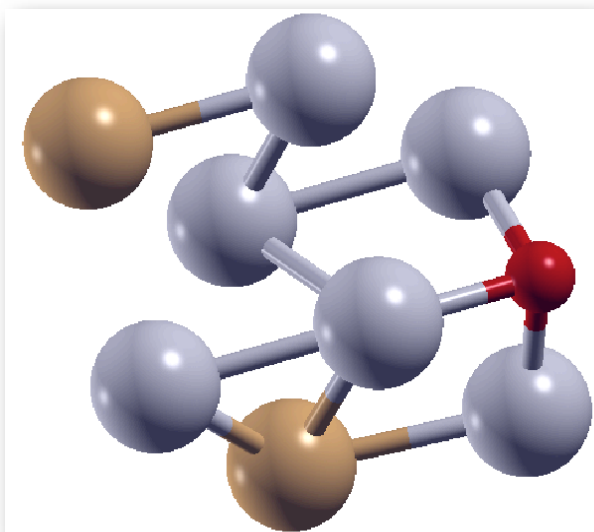
# Comparison of other charge mixing schemes



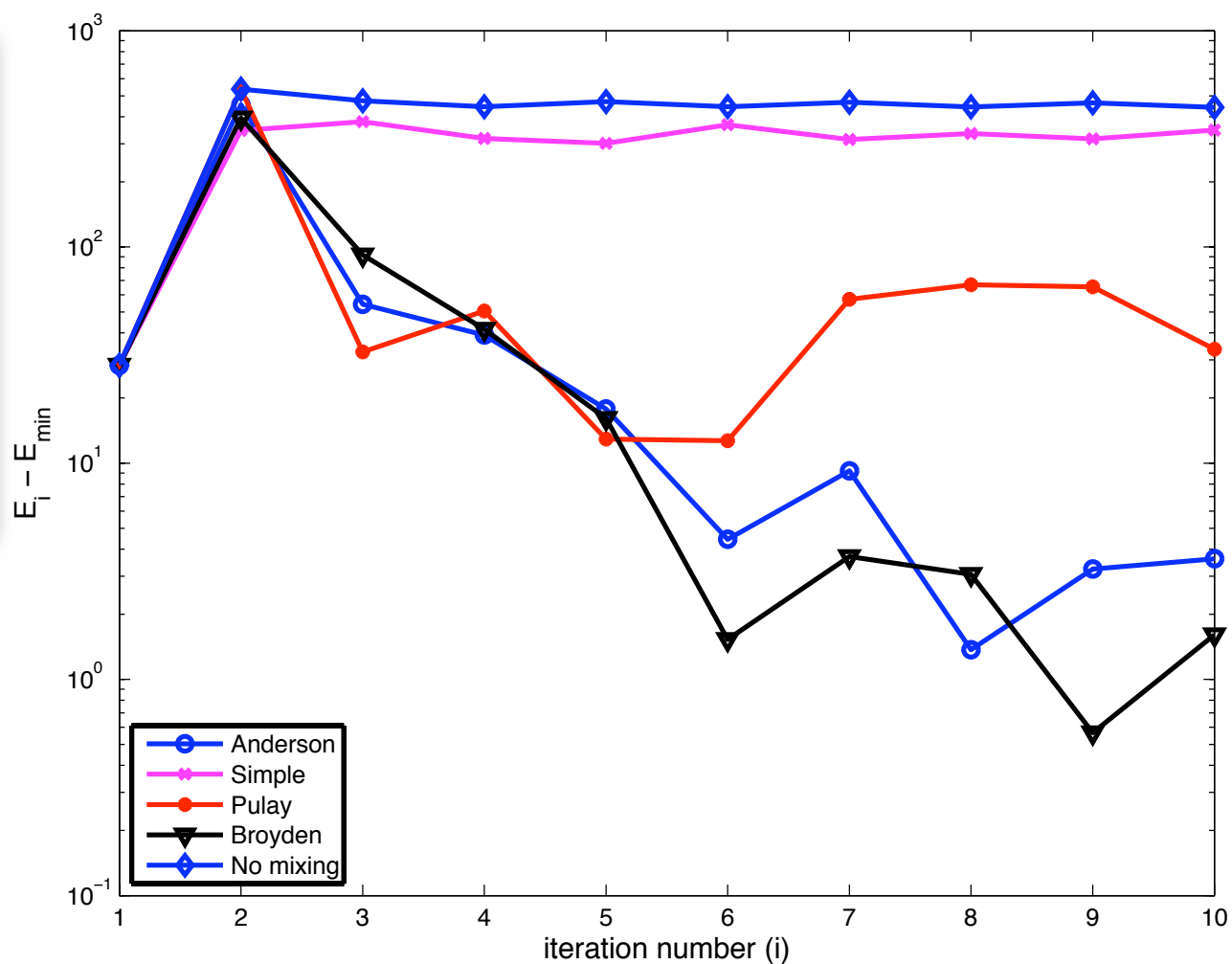
Alanine



# Charge mixing can fail



Pt<sub>6</sub>Ni<sub>2</sub>O



# Direct Constrained Minimization

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- Assume  $x^{(i)}$  is the current approximation
- Idea: minimize the energy in a certain (smaller) subspace
- Update  $x^{(i+1)} = \alpha x^{(i)} + \beta p^{(i-1)} + \gamma r^{(i)}$ ;
  - $p^{(i-1)}$  previous search direction;
  - $r^{(i)} = H^{(i)} x^{(i)} - \theta^{(i)} x^{(i)}$ ;
  - choose  $\alpha$ ,  $\beta$  and  $\gamma$  so that
    - \*  $x_{k+1}^T x_{k+1} = 1$ ;
    - \*  $E(x_{k+1}) < E(x_k)$ ;

*Remark 1: A nonlinear CG-like algorithm*

*Remark 2: Extension of LOBPCG (Knyazev) to nonlinear EV*

# Subspace minimization

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- Let  $V = (x^{(i)}, p^{(i-1)}, r^{(i)})$ ;  $x^{(i+1)} = Vy$ , for some  $y$ ;

- Solve

$$\min_{y^T V^T V y = 1} E(Vy)$$

- Equivalent to solving

$$\begin{aligned} G(y)y &= \lambda B y \\ y^T B y &= 1 \end{aligned}$$

where  $B = V^T V$  and  $G(y) = V^T [L + \alpha \text{Diag}(L^{-1} \rho(Vy))] V$



# DCM Algorithm

- Input: Initial guess
- Output:  $X$  such that  $E_{KS}$  is minimized
  1.  $P^{(0)} = \emptyset, i = 0;$
  2. while ( not converged )
    - (a)  $\Theta^{(i)} = X^{(i)*} H^{(i)} X^{(i)};$
    - (b)  $R^{(i)} = H^{(i)} X^{(i)} - X^{(i)} \Theta^{(i)};$
    - (c) Set  $Y = (X^{(i)}, P^{(i-1)}, K^{-1} R^{(i)})$ ;
    - (d) Solve  $\min_{G^* Y^* Y G = I_k} E_{tot}(Y G);$
    - (e)  $X^{(i+1)} = Y G(1 : n_e, :); P^{(i+1)} = Y G(n_e + 1 : 3n_e, :);$
    - (f)  $i \leftarrow i + 1;$

*C. Yang, J. Meza, L. Wang, A Constrained Optimization Algorithm for Total Energy Minimization in Electronic Structure Calculation, J. Comp. Phys., 217 709-721 (2006)*

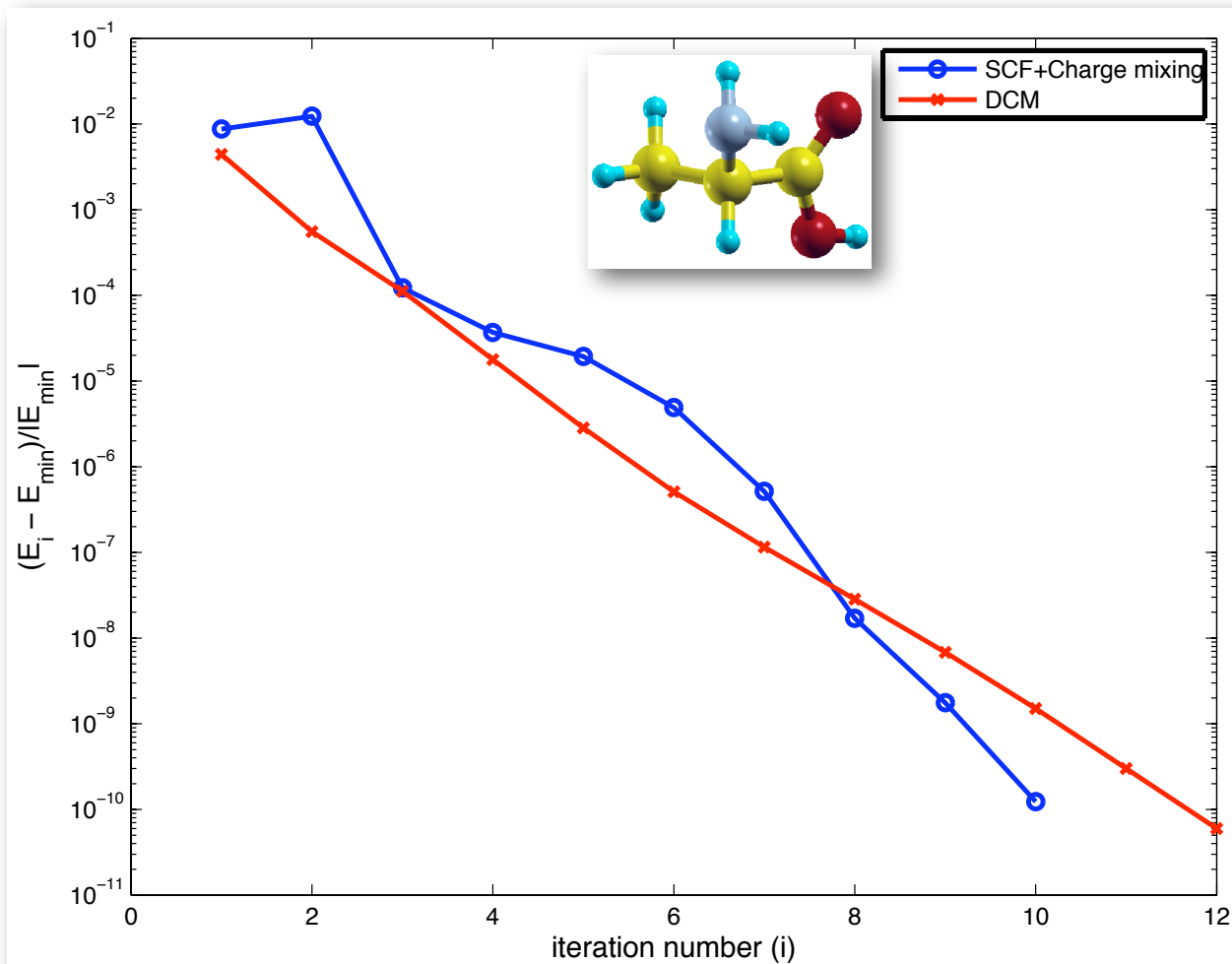
# Test problems

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- KSSOLV Matlab code for solving the Kohn-Sham equations
  - Open source package
  - Handles SCF, DCM, Trust Region
  - Various mixing strategies
- Example problems: alanine and graphene
- Tests run on desktop computer

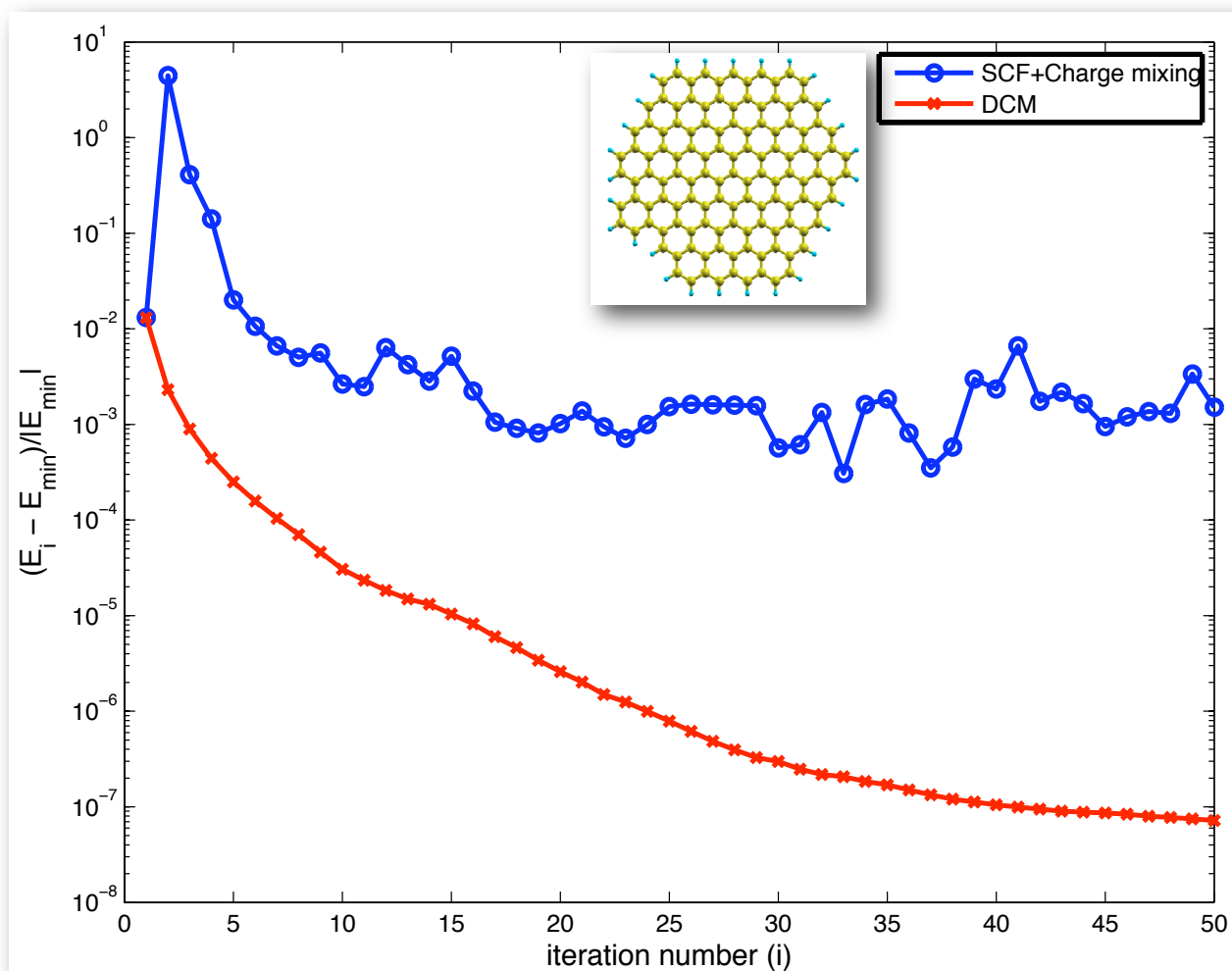
# Example 1: Alanine

- sampling grid:
  - 96 x 48 x 96  
(ecut=25 Ryd)
- 10 PCG iterations / SCF outer iteration
- 3 inner SCF iteration / DCM outer iteration
- supercell:
  - 20 x 15 x 20



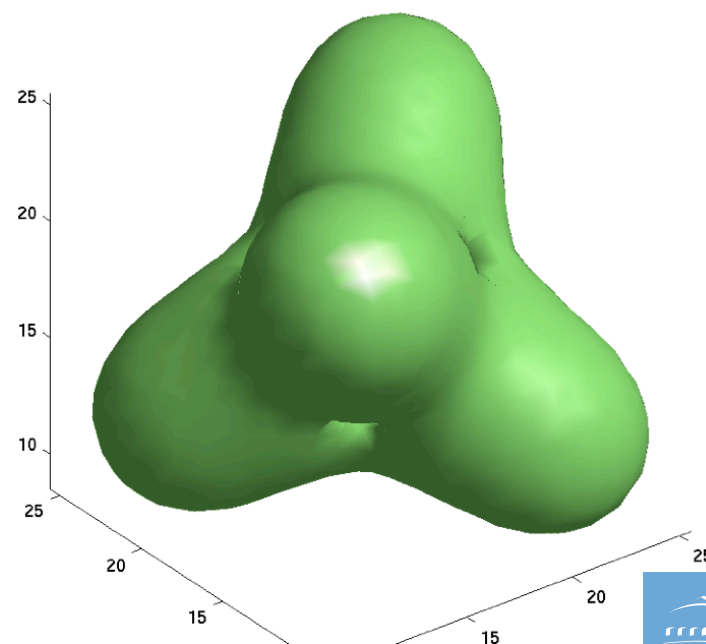
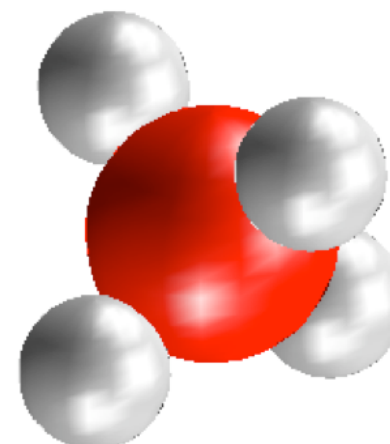
## Example 2: Graphene

- sampling grid:
  - 114 x 114 x 15
- 10 PCG iterations / SCF outer iteration
- 5 inner SCF iteration / DCM outer iteration
- supercell:
  - 40 x 40 x 5



# Example: SiH<sub>4</sub>

```
a1 = Atom('Si');
a2 = Atom('H');
alist = [a1 a2 a2 a2 a2];
xyzlist= [
  0.0    0.0    0.0
  1.61   1.61   1.61
  ... ];
mol = Molecule();
mol = set(mol,'Blattice',BL);
mol = set(mol,'atomlist',alist);
mol = set(mol,'xyzlist',xyzlist);
mol = set(mol,'ecut', 25);
mol = set(mol,'name','SiH4');
...
[Etot, X, vtot, rho] = dcm(mol);
isosurface(rho);
```



# Comparison of DCM vs. SCF using KSSOLV

system	SCF time	DCM time	SCF error	DCM error
$C_2H_6$	26	25	9.4 e-6	3.5 e-6
$CO_2$	26	23	3.1 e-3	1.1 e-4
$H_2O$	16	16	5.7 e-5	2.2 e-5
$HNCO$	34	32	7.4 e-3	6.8 e-5
Quantum dot	18	16	5.0 e-3	3.7 e-1
$Si_2H_4$	25	23	1.8 e-3	2.7 e-4
silicon bulk	15	15	3.0 e-4	9.6 e-6
$SiH_4$	20	19	9.7 e-6	4.9 e-7
$Pt_2Ni_6O$	415	281	3.7 e0	4.9 e-2
pentacene	887	493	5.2 e-1	2.5 e-2

# Summary

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- Despite dire warnings, mathematical techniques actually help in chemistry
- New approach for solving the Kohn-Sham equations using a direct optimization method improves convergence
- Trust region modification increases robustness of both SCF and DCM
- New computational software tools for modeling and simulation of nanosystems

# Where do we go from here

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- Investigate new algorithms to speed up convergence
- Develop more accurate methods
- Expand applicability of methods to new systems, (metals)
- Develop linear scaling versions of DCM



# First Nanoscientists?

The New York Times

February 21, 2005

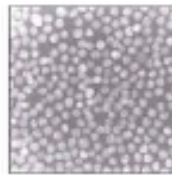


## The First Nanotechnologists

Ancient stained-glass makers knew that by putting varying, tiny amounts of gold and silver in the glass, they could produce the red and yellow found in stained-glass windows. Similarly, today's scientists and engineers have found that it takes only small amounts of a nanoparticle, precisely placed, to change a material's physical properties.

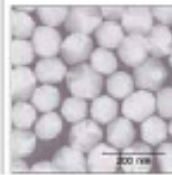
### Gold particles in glass

Size: 25 nm  
Shape: sphere  
Color reflected:  
  
100 nanometers is  
0.0001 millimeter



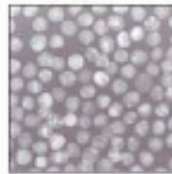
### Silver particles in glass


Size: 100 nm  
Shape: sphere  
Color reflected:  

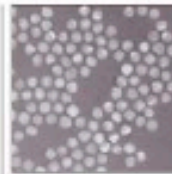



Had medieval artists been able to control the size and shape of the nanoparticles, they would have been able to use the two metals to produce other colors. Examples:

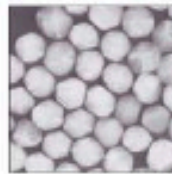
Size: 50 nm  
Shape: sphere  
Color reflected:  

Size: 40 nm  
Shape: sphere  
Color reflected:  




Size: 100 nm  
Shape: sphere  
Color reflected:  

Size: 100 nm  
Shape: prism  
Color reflected:  




Source: Dr. Chad A. Mirin, Institute of Nanotechnology, Northwestern University

\*Approximate

